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HYPERON POLARIZATION NEWS

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Fermilab Hyperon '99 Symposium

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NATURE HAS NO DIFFICULTY PRODUCING POLARIZED HYPERONS

Brief History

1.) $p \rightarrow \Lambda$ $\vec{P}_\Lambda$ opposite to $\vec{k}_{in} \times \vec{k}_{out}$

SU(6) baryon wave functions + 5 quarks with
spin down first proposed in our group
by Ken Heller

$\vec{P}_\Lambda = \vec{P}_S$ gives correct sign for
 $P_{\Sigma^+}, P_{\Sigma^-}, P_{\Sigma^0}, P_{\Xi^-}$

predicts $\vec{P}_{\Sigma^+} = -\frac{1}{3} \vec{P}_\Lambda$ too small

2.) $\vec{P}_{\Sigma^+} \sim -\vec{P}_\Lambda$ so need polarized diquarks too

$$\left. \begin{array}{l} \text{single quark pol } \uparrow \\ \text{diquark pol } \downarrow \end{array} \right\} \text{DeGrand + Miettinen}$$

DGM MODEL CONT'D

LEADING PARTONS DECELERATE AND ACQUIRE POS POL
TRAILING PARTONS ACCELERATE AND ACQUIRE NEG POL

Polarizations for various transitions predicted by the leading-partons-trailing-partons model.

$B \rightarrow B'$ transition	Polarization
$p \Rightarrow n, \Sigma^+ \Rightarrow \Xi^0,$ $\Sigma^- \Rightarrow \Xi^-, \Sigma^\pm \Rightarrow \Sigma^0,$ $\Xi^0 \Rightarrow \Sigma^0$	$-\frac{20}{21}\epsilon + \frac{1}{42}\delta$
$p \Rightarrow \Sigma^+, n \Rightarrow \Sigma^-,$ $p \Rightarrow \Sigma^0, n \Rightarrow \Sigma^0,$ $\Xi^0 \Rightarrow \Xi^-$	$\frac{1}{3}\epsilon + \frac{2}{3}\delta$
$p, n \Rightarrow \Lambda$	$-\epsilon$
$\Sigma^\pm \Rightarrow \Lambda, \Xi^0, \Xi^- \Rightarrow \Lambda$	$-\frac{2}{3}\epsilon + \frac{1}{6}\delta$
$p \rightarrow \Sigma^-, n \rightarrow \Sigma^+,$ $\Sigma^+ \rightarrow \Xi^-, \Sigma^- \rightarrow \Xi^0,$ $\Xi^0 \rightarrow p, \Xi^- \rightarrow n$	$\frac{2}{3}\epsilon - \frac{1}{6}\delta$
$p \Rightarrow \Xi^-, n \Rightarrow \Xi^0$ $\Sigma^+ \Rightarrow \Sigma^-, p \rightarrow \Xi^0,$ $\Xi^0 \rightarrow \Sigma^-, \Sigma^- \rightarrow p, n \rightarrow \Xi^-$ $\Xi^- \rightarrow \Sigma^+, \Sigma^+ \rightarrow n$	$-\frac{1}{3}\epsilon - \frac{2}{3}\delta$
$\pi, K^+ \rightarrow \Lambda$	$-\delta/2$
$K^- \rightarrow \Lambda$	ϵ

EXPTL EVIDENCE $\Rightarrow \epsilon \sim \delta$ TO 20%

DG+M SCORE CARD

1985

TABLE I. Comparison of the model of Ref. 11, with $\epsilon=6$, with data from unpolarized beams.

Transition	Predicted asymmetry	Observed asymmetry	Energy (GeV)	Reference
$p \rightarrow \Lambda$	$-\epsilon$	-0.1 to -0.2	24-2000	1
$p \rightarrow \bar{\Lambda}$	0	0	24-2000	1
$p \rightarrow \Sigma^+$	ϵ	0.1 to 0.2	400	2
$p \rightarrow \Sigma^-$	$\epsilon/2$	0.15, 0.3	400	3,4
$p \rightarrow \Sigma^0$	ϵ			
$p \rightarrow \Xi^0$	$-\epsilon$	-0.1 to -0.2	400	5
$p \rightarrow \Xi^-$	$-\epsilon$	-0.1 to -0.2	400	6
$K^+ \rightarrow \bar{\Lambda}$	ϵ	$> 0.4, x_F > 0.3$	32,70	7
$K^- \rightarrow \Lambda$	ϵ	0.4	14	8
$\pi^- \rightarrow \Lambda$	$-\epsilon/2$	-0.05	18	9
$\gamma \rightarrow \Lambda$	$-\epsilon/2$	-0.1	20	10

WAGG $\Sigma^- \rightarrow \Xi^-$ - 6 OK WITHIN ERRORS 330 GeV

$x_F > 0$
 $p_T > 1$ { ACCMOR $\pi^- \rightarrow \Lambda$ - $\epsilon/2$ - 0.28 ± 0.09 230 GeV
 " $\pi^- \rightarrow \bar{\Lambda}$ - 0.01 ± 0.12 "

ALMOST EQUAL NUMBERS of Λ & $\bar{\Lambda}$

DG+M SCORE CARD CONT'D

POLARIZATION BY OTHER BARYONS

WA 89 $\Sigma^- \rightarrow \Sigma^+$ - e -0.035 ± 0.012 $\langle P \rangle = 10$
 $\langle X \rangle = 0.27$
 3306 eV

OK, BUT SMALL X

SELEX $\Sigma^- \rightarrow \Lambda$ - e/2 SEE PLOT -
 WA 89 WRONG SIGN

WA 89 DATA CONSISTENT WITH SELEX,
 BUT LED TO OPPOSITE CONCLUSION.

NEUTRAL BEAM PRODUCTION 1.8 MRAD 3956 eV NEG
 BEAM

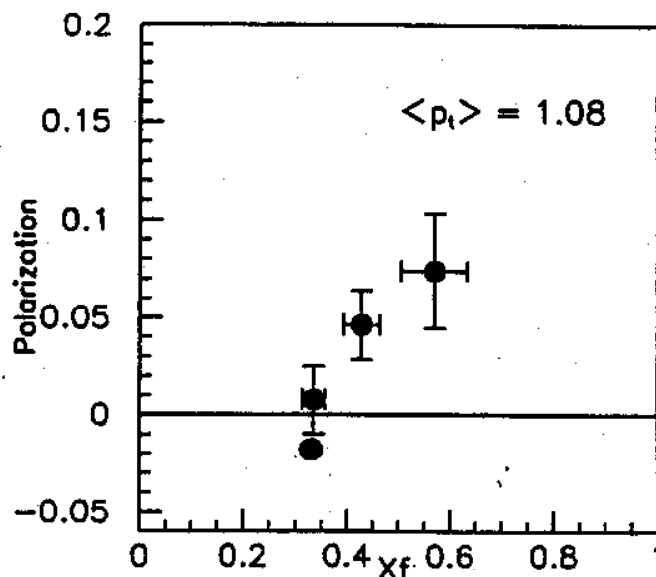
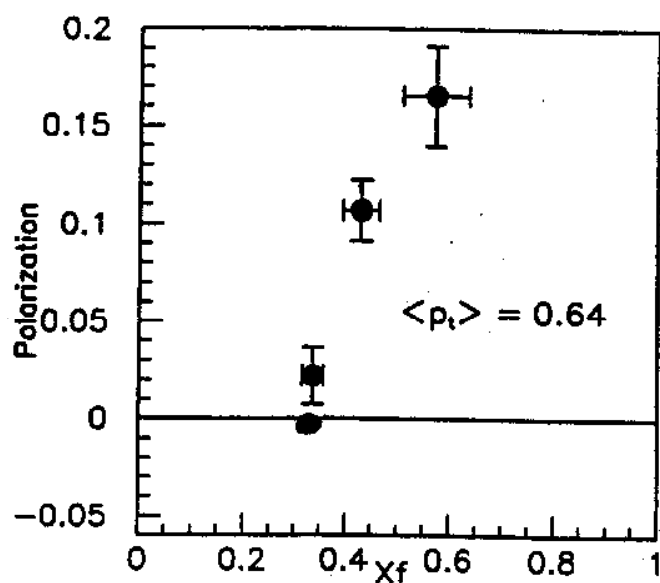
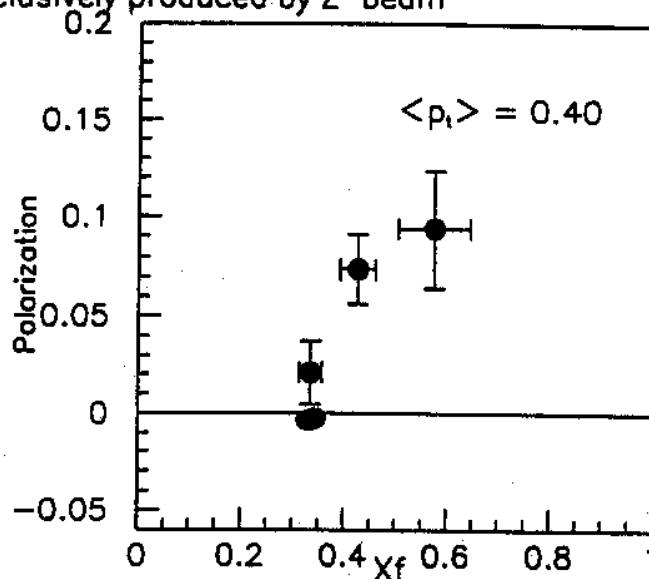
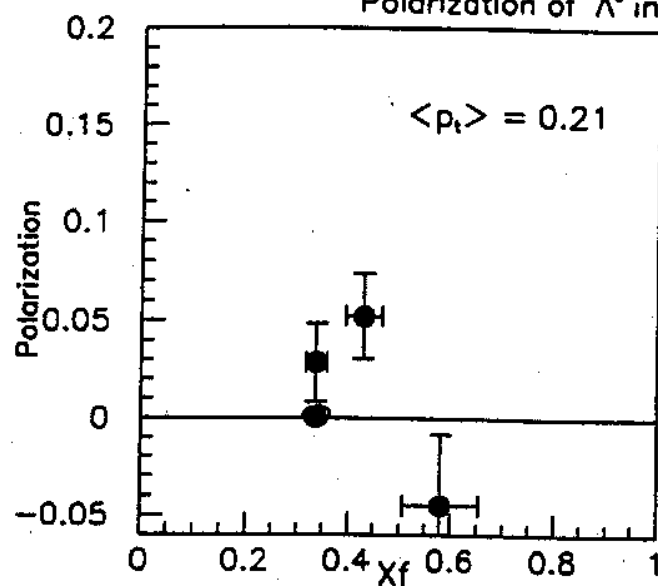
$(n, \Lambda, \Xi^0)_{\text{unpol}} \rightarrow \Xi^-$ $P_{\Xi^-} = 0.002 \pm 0.0042$

$(n, \Lambda, \Xi^0)_{\text{unpol}} \rightarrow \Sigma^-$ $P_{\Sigma^-} = +0.053 \pm 0.012$

$p_{\text{unpol}} \rightarrow \Sigma^-$ $P_{\Sigma^-} = -0.01 \pm 0.01$

SELEX COLLAB UNPUBLISHED DATA

● WA89 330 GeV Σ^-
 Z Phys A 350, 379 (1980)

Polarization of Λ^0 inclusively produced by Σ^- beam

WA89

 $\langle p_t \rangle \Lambda p.1$ 1.32 -0.055 ± 0.005 1.80 -0.033 ± 0.020

$$A = k_{beam}^{\rightarrow} \times k_{part}^{\rightarrow}$$

DG + M SCORE CARD CONT'D

SPIN EXCHANGE AND Λ ASYMMETRY

FOR $p \uparrow \rightarrow \Lambda$ 200 GeV p 's E704

THE Λ SHOULD NOT KNOW ANYTHING
ABOUT THE PROTON SPIN

BUT IT DOES.

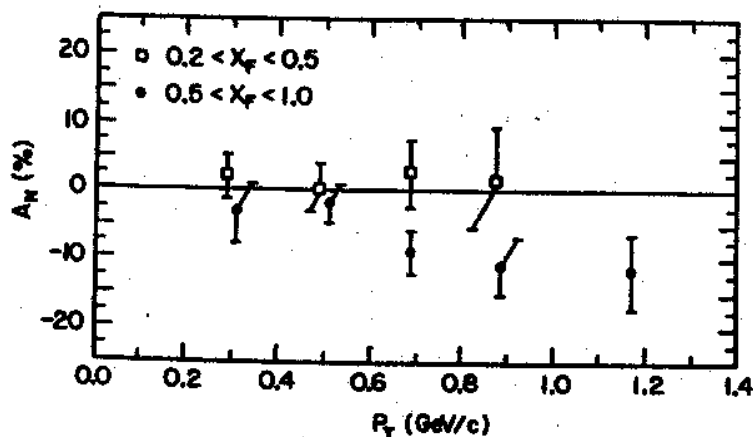


FIG. 4. A_N data for $p \uparrow + p \rightarrow \Lambda^0 + X$ as a function of p_T .

Λ ASYMMETRY RESEMBLES THE Λ
POLARIZATION, JUST LIKE ELASTIC SCATTERING!

SPIN EXCHANGE

NEG POLARIZED NEUTRAL BEAM $\rightarrow$ 3956eV
NEG BEAM

$$(n, \Lambda, \Xi^0) \rightarrow \Xi^- \quad P_{\Xi^-} = -0.1172 \pm 0.0062$$

$$(n, \Lambda, \Xi^0) \rightarrow \Sigma^- \quad P_{\Sigma^-} = -0.076 \pm 0.021$$

THESE RESULTS MAKE SENSE.

ANTI HYPERON POLARIZATION

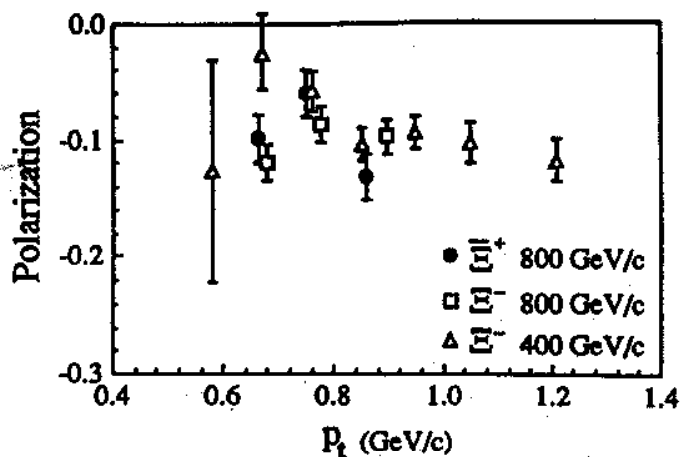
$$p \rightarrow \bar{\Lambda} \quad \text{NO POLARIZATION}$$

$$p \rightarrow \bar{\Xi}^+ \quad P \sim P_{\Xi^-} \neq 0$$

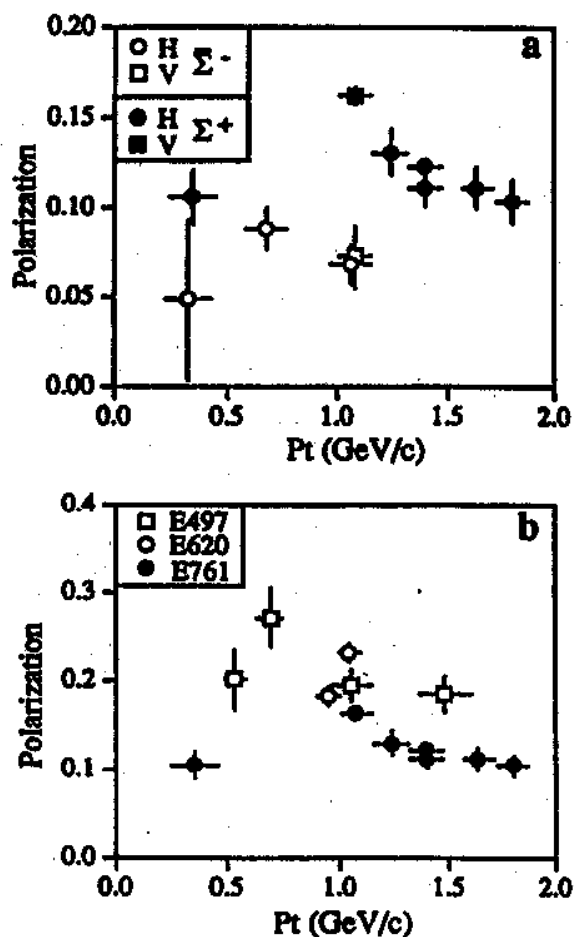
$$p \rightarrow \bar{\Sigma}^- \quad P \sim P_{\Sigma^+} \neq 0$$

$$p \rightarrow \bar{\Xi}^0 \quad \text{NO POLARIZATION!} \quad KTeV$$

Ho, et al



MORELOS, et al



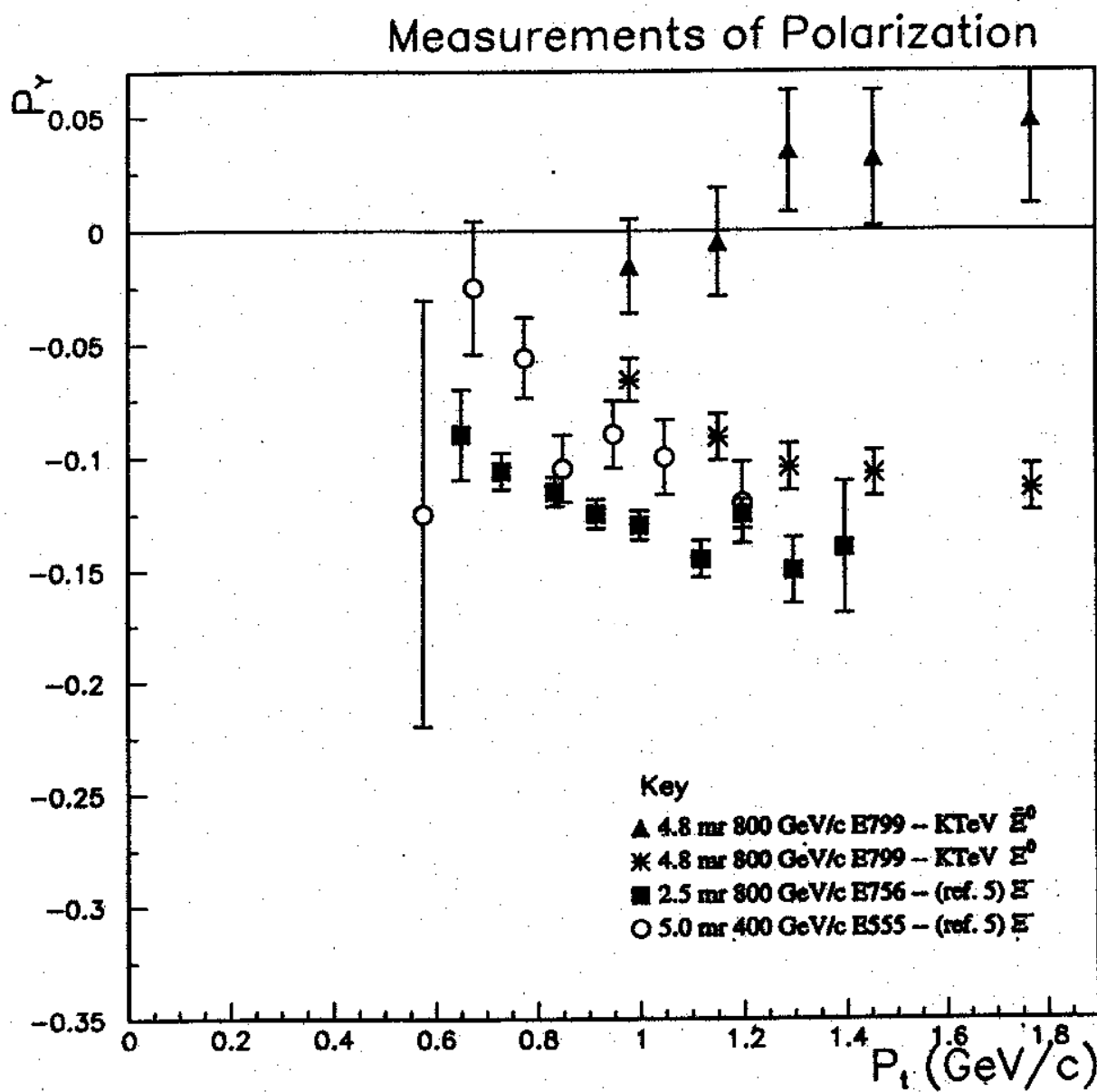
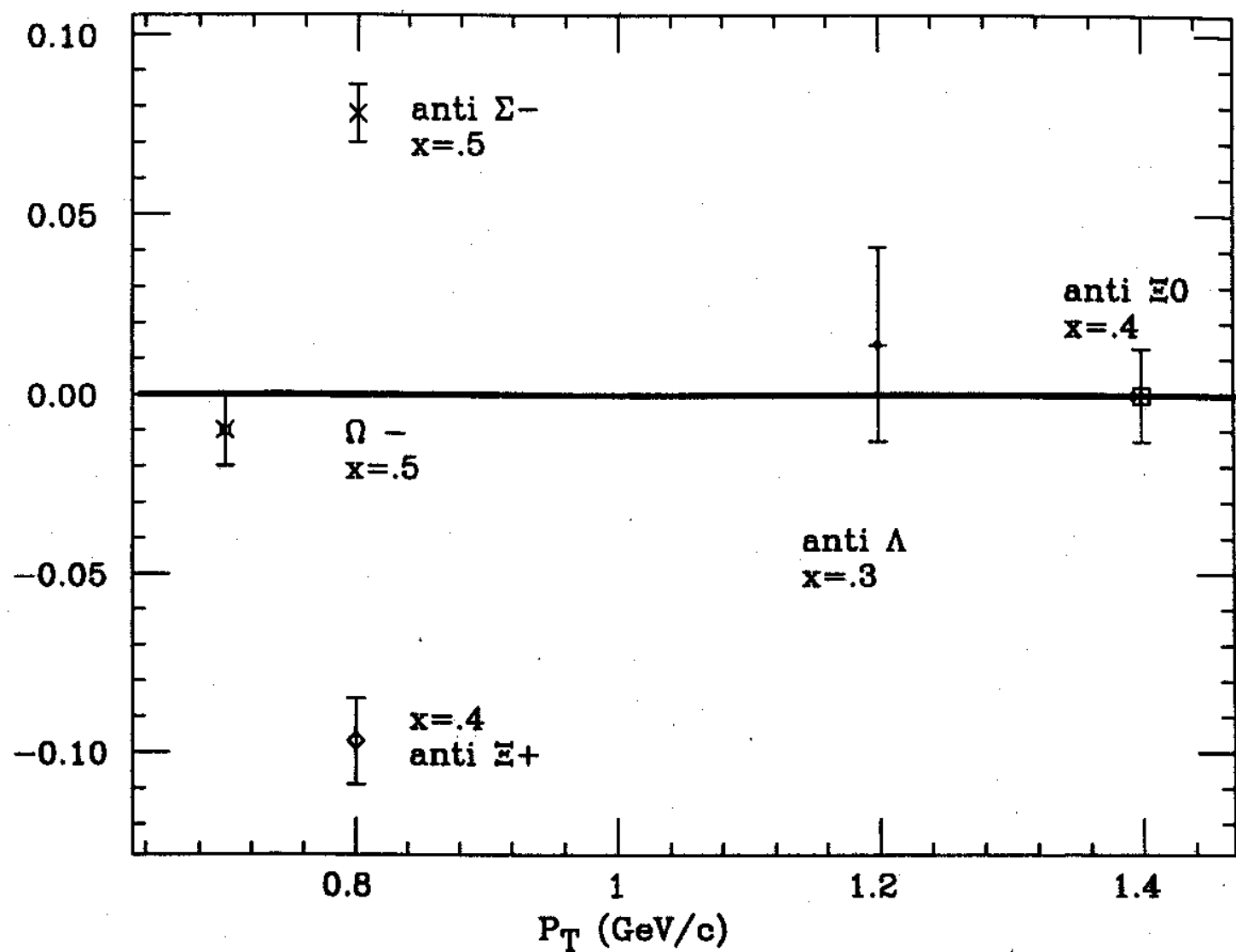


Figure 7.8: More past results plotted against transverse momentum.

anti hyperon polarization, 800 GeV protons



THE PION IDEA

SOME AUTHORS, like K. Kubo et al have
SUGGESTED THAT $\bar{\Upsilon}$ POLARIZATION COULD
BE EXPLAINED BY FAST PIONS

$$\left. \begin{array}{l} p \rightarrow n \pi^+ \\ p \rightarrow \Delta^{++} \pi^- \end{array} \right\} \text{pion} \rightarrow \Upsilon \bar{\Upsilon}$$

REMEMBER $\pi^- \rightarrow \Lambda, \bar{\Lambda}$ IN APPX EQUAL NUMBERS

AND $P_{\Lambda} = 0$, CONSISTENT WITH $p \rightarrow \bar{\Lambda}$

HAS $\pi \rightarrow \Sigma, \bar{\Sigma}$ BEEN MEASURED?

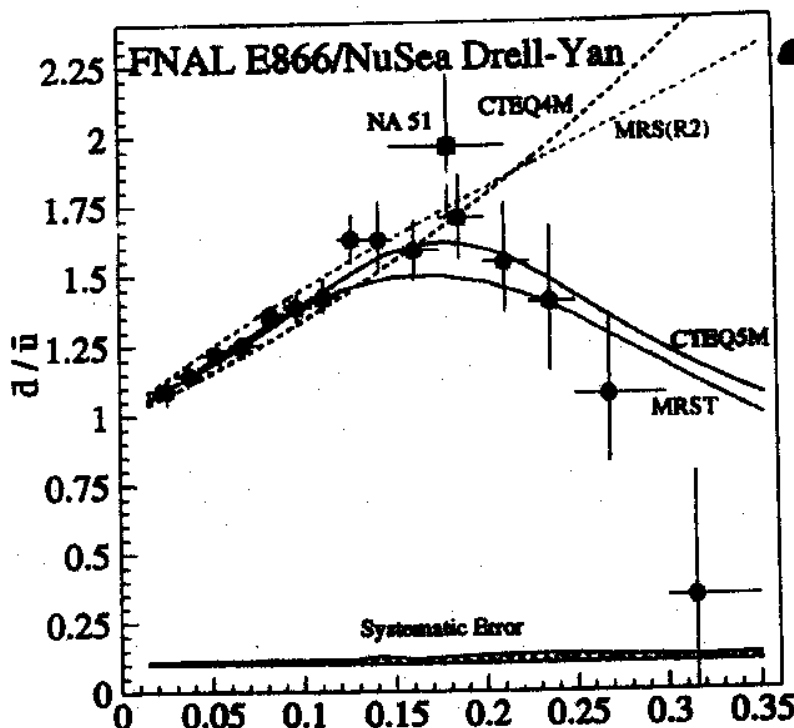
THESE IDEAS, like OPE, ARE ENJOYING
A RENAISSANCE

FOR INSTANCE $\frac{\bar{\sigma}}{\bar{n}} > 1$

MAX KLEIN LP'99

15

flavour asymmetry in the nucleon sea $\bar{u} \neq \bar{d}$

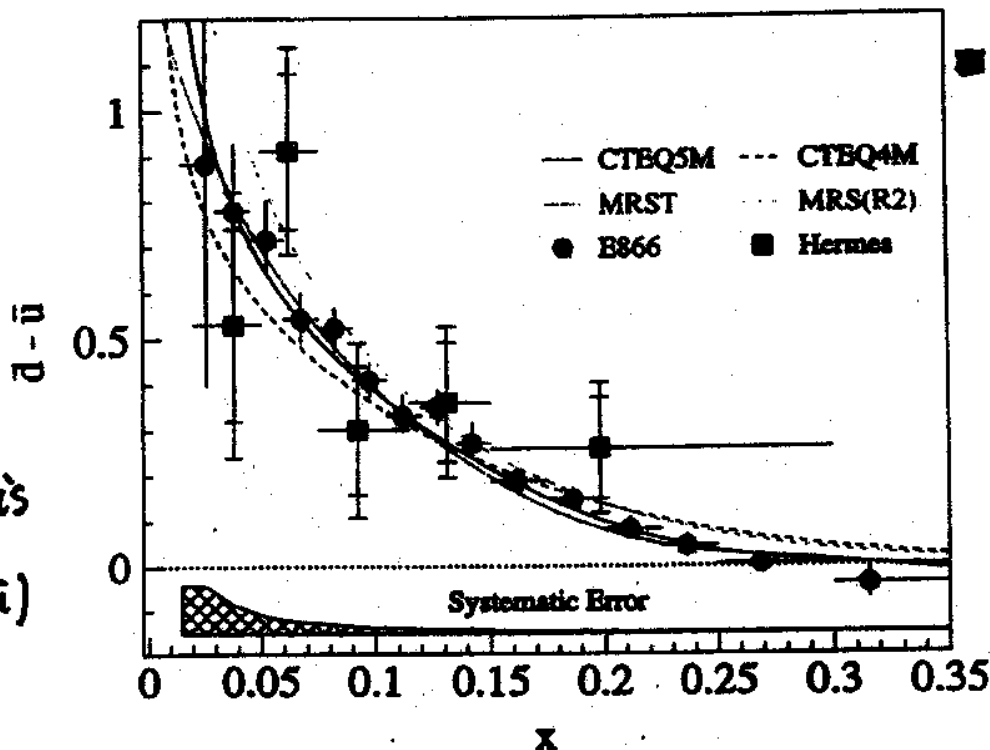


● E866/NuSea
pp, pd DY.
 $3.7 \cdot 10^5 \mu\text{pairs}$
 $\langle Q^2 \rangle \approx 54 \text{ GeV}^2$

$$\int_0^1 (\bar{u} - \bar{d}) dx$$

Gottfried SR: $\int \frac{1}{x} (F_2^p - F_2^n) dx = \frac{1}{3} + \frac{2}{3} \int (\bar{u} - \bar{d}) dx$

E86: -0.10 ± 0.04
NMC: -0.15 ± 0.04



■ HERMES
 $eN \rightarrow e\bar{\nu}^+ X$
p, d
 $\langle Q^2 \rangle = 2.3 \text{ GeV}^2$

↑
this is
not
 $x(\bar{d} - \bar{u})$

OTHER NEW RESULTS

1.) KTeV $\phi \rightarrow \Xi^0$ 800 GeV p'

2.) E781 $\phi \rightarrow \Sigma^+$ 800 GeV p'

3) EXCLUSIVE CHANNELS AT BNL

27.5 GeV PROTONS

$\phi\phi \rightarrow p \Lambda K^+ \pi^+ \pi^-$

EVENTS

5421

$\pi^+ \pi^- \pi^+ \pi^-$

51195

$\pi^+ \pi^- \pi^+ \pi^- \pi^+ \pi^-$

48195

8 CHARGED π 's

14582

QUOTE ABOVE FIT ADEQUATE χ^2

$$P_{\Lambda} = (-0.443 \pm 0.032) \chi_p \beta_T \quad -1 \leq \chi_F \leq +1$$

$$0.5 \beta_T \leq 1.860V$$

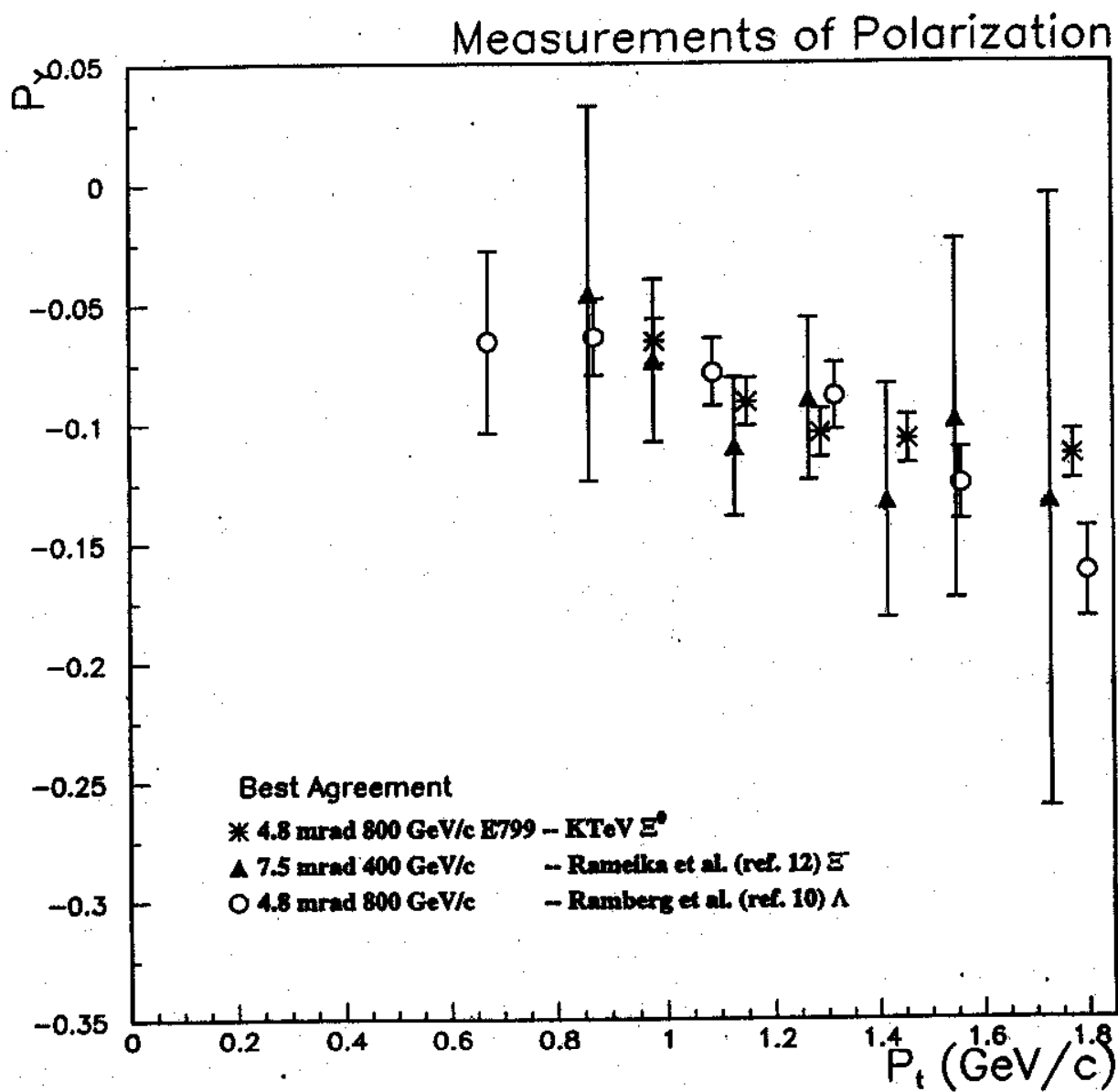


Figure 7.9: Past results plotted with ours versus transverse momentum.

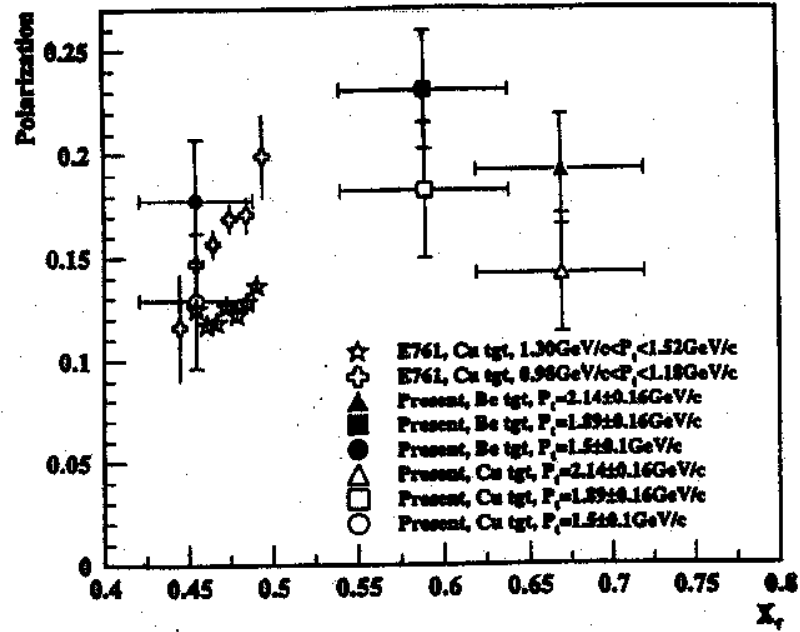


Figure 4: X_T dependence of the Σ^+ polarisation (only statistical errors are shown). Open symbols correspond to the Cu production target, filled symbols - to the Be target. Open crosses and stars represent the E761 data [2] at $p_t \approx 1$ GeV/c and $p_t \approx 1.5$ GeV/c, respectively. Circles correspond to the present measurement at $p_t \approx 1.50$ GeV/c. Boxes correspond to the present measurement at $p_t \approx 1.89$ GeV/c. And, finally, triangles correspond to the present measurement at $p_t \approx 2.14$ GeV/c.

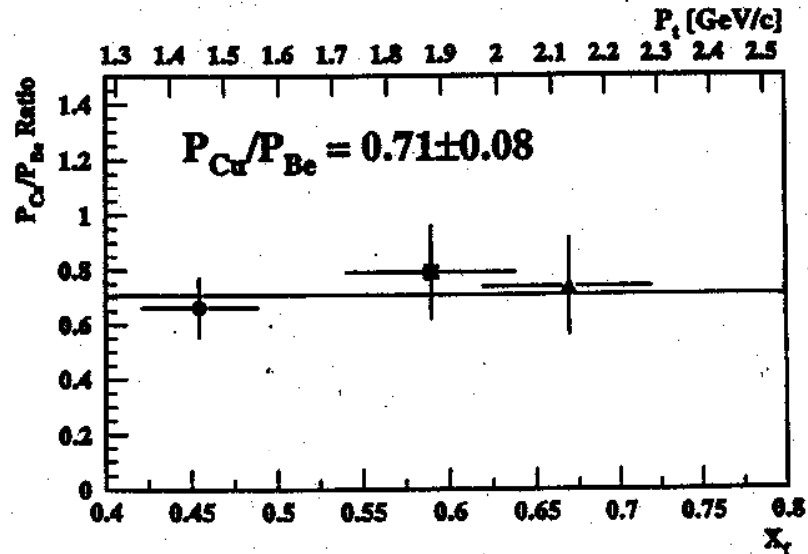
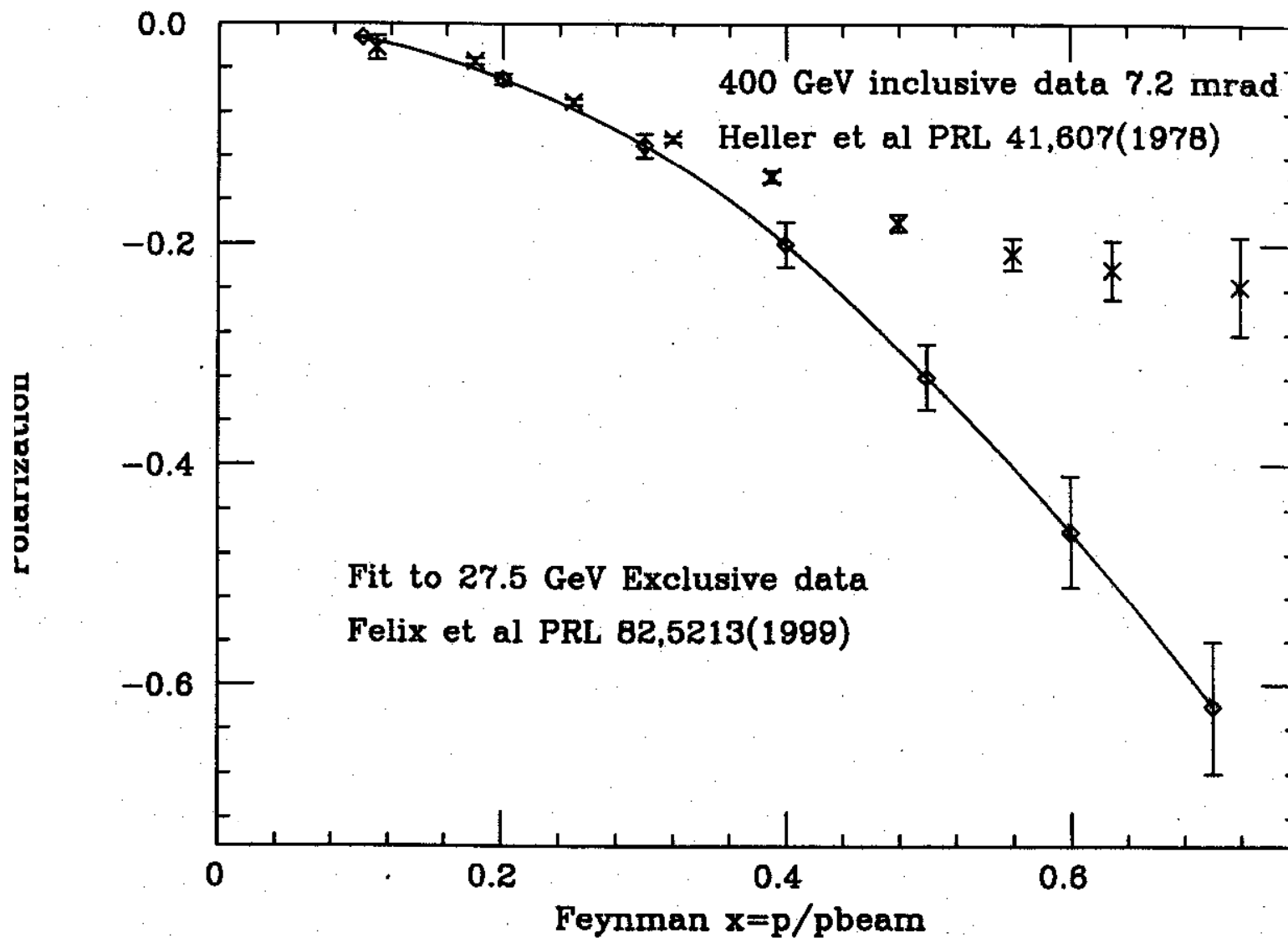


Figure 5: Ratio of the Σ^+ polarization values for Cu and Be targets (only statistical errors are taken into account). The horizontal line corresponds to the average ratio of 0.71 ± 0.08 .

Polarization of Λ hyperons by protons



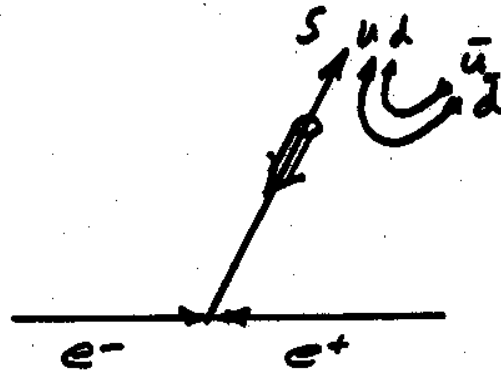
DECELERATING STRANGE QUARKS

AND

$$Z \rightarrow s\bar{s}$$

LARGE LONGITUDINAL POLARIZATION FOR STRANGE QUARKS

$S \rightarrow \Lambda$ VIA DIQUARK HADRONIZATION



$u d$ SINGLET SPIN STATE

SOURCES of Λ in $Z \rightarrow q \bar{q}$

① LEADING S QUARK $\rightarrow \Lambda$ $P_\Lambda = P_S$

② LEADING S QUARK $\rightarrow \Sigma, \Sigma^*, \Xi, \Xi^*$ $P_\Lambda \neq 0$

③ LEADING u/d QUARK $\rightarrow \Lambda$

④ Λ CONTAINS NO LEADING QUARKS

$\left. \begin{array}{l} \text{③} \\ \text{④} \end{array} \right\} P_\Lambda = 0$
LONGITUDINAL

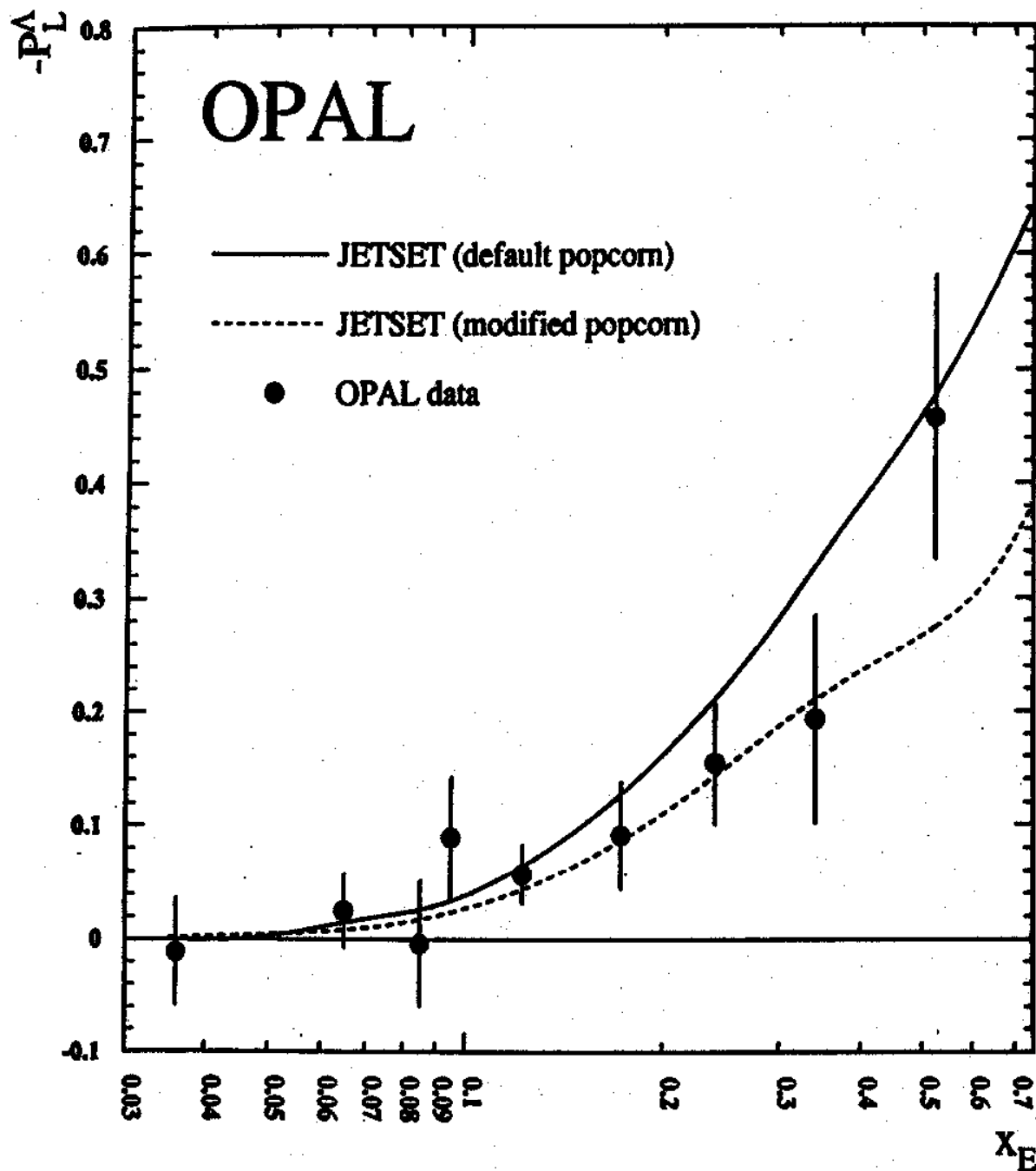
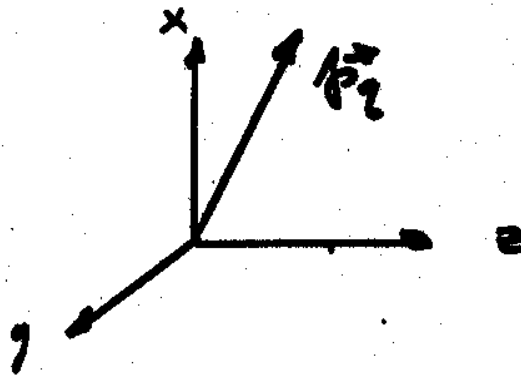


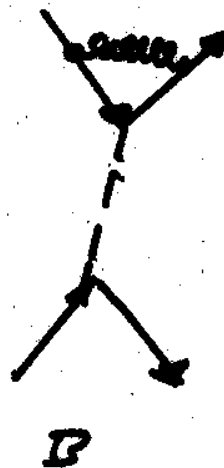
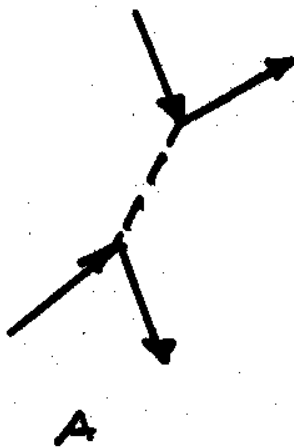
Figure 3: The longitudinal polarization predicted using tuned versions of JETSET with different versions of the popcorn model of baryon production. The measurements are shown as solid points and the error bars are the statistical plus systematic errors, added in quadrature.

TRANSVERSE POLARIZATION IN $Z \rightarrow g\bar{g}$

NORMAL TO $e^+e^- \rightarrow g\bar{g}$ PLANE $\hat{n} = \hat{y}$



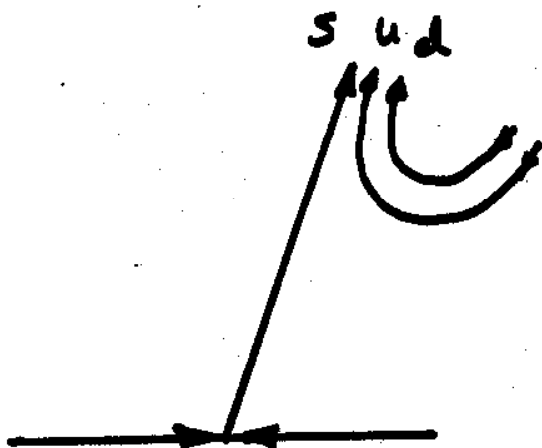
VANISHES IN LOWEST ORDER (A)



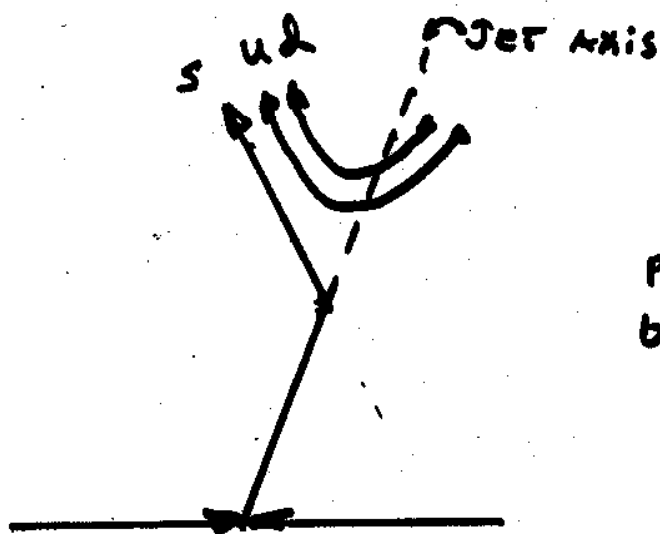
A+B interference $\rightarrow \vec{P} \cdot \hat{n} = \frac{4}{3} \alpha_s \frac{M_g}{Q} \frac{\sin\theta \cos\theta}{1 + \cos^2\theta}$

Kane, Pomplun, & Repko

POLARIZATION IN FRAGMENTATION



TRANSVERSE TO WHAT? Well, how ABOUT



PICTURE SIMILAR TO
BEAM FRAGMENTATION
MESON $\rightarrow \Lambda$?

SO CORRELATE THE Λ SPIN WITH THE NORMAL
TO THE JET- Λ PLANE.

S. Gaur/Layer CTA/ PRL 56 2244 (1986)
 OPPOSITE SIGN CONVENTION

$$\bar{M} = \frac{1}{p_A} \times \frac{p_{beam}}{p_A \times p_B}$$

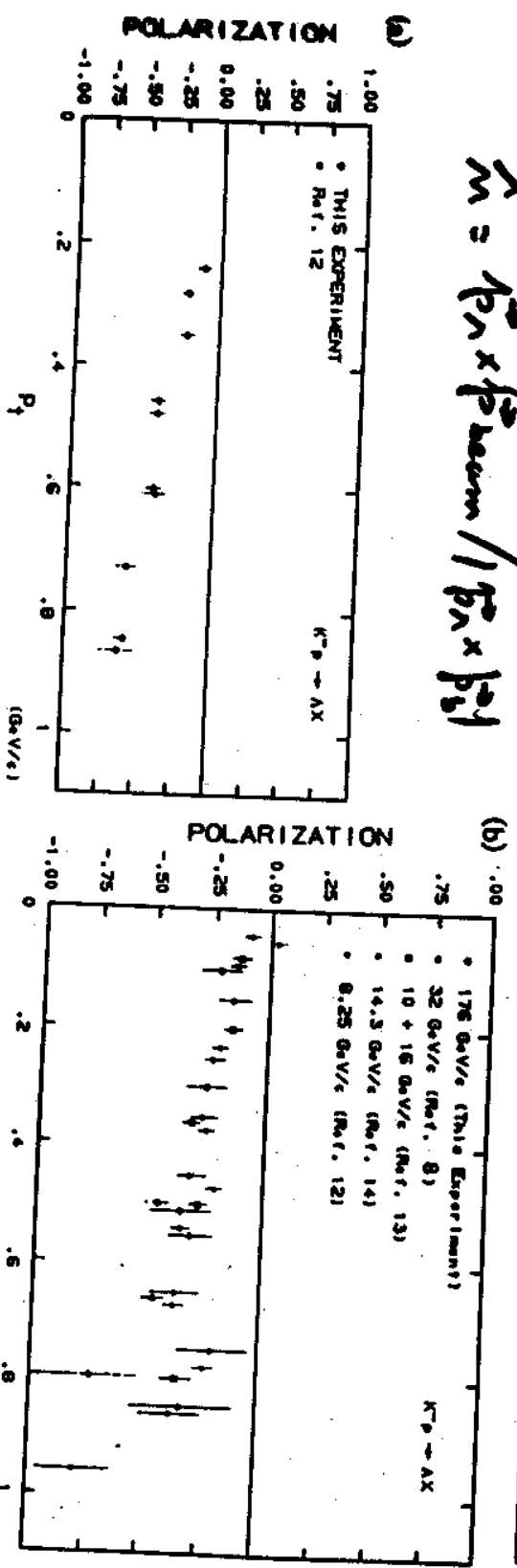


FIG. 5. (a) Lambda polarization from $K^- + p \rightarrow \Lambda + X$ and the 8.25-GeV/c data of Ref. 12 plotted vs p_T . (b) Lambda polarization from $K^- + p \rightarrow \Lambda + X$ and the data from Refs. 13 and 14 plotted vs X_F .

$K^+ \rightarrow \bar{\Lambda}$ SIMILAR - SAME SIGN
 Ahinen et al PRL 52, 183 (1983).

USING THE CONVENTION $\hat{n} = \frac{\vec{p}_{in} \times \vec{p}_{out}}{|\vec{p}_{in} \times \vec{p}_{out}|}$

$K^- p \rightarrow \Lambda X$ $\vec{P}_\Lambda \cdot \hat{n}$ positive

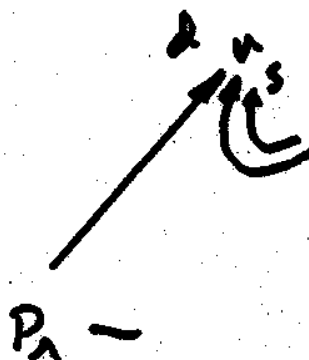
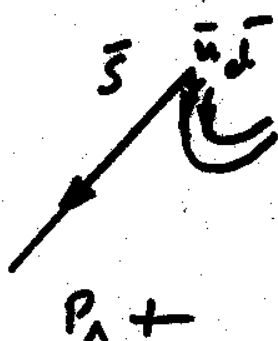
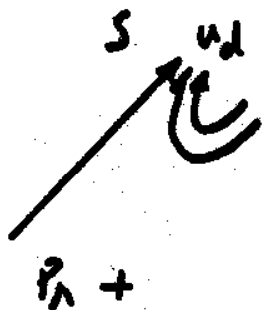
$K^+ p \rightarrow \bar{\Lambda} X$ $\vec{P}_\Lambda \cdot \hat{n}$ positive

EACH $P_\Lambda \sim 20\%$ IN BEAM FRAGMENTATION REGION

$\pi^- Cu \rightarrow \Lambda$ $x_F > 0$ $p_T > 1.6 \text{ GeV}$

$P_\Lambda = -0.28 \pm 0.09$ OPPOSITE SIGN

ACCMOR COLLAB PL 3258, 531 (1984)



EXPERIMENTAL DATA

$$Z \rightarrow \Lambda, \bar{\Lambda} X$$

$$\vec{P}_\Lambda \cdot \hat{n}$$

$$\hat{n} = \vec{p}_\Lambda \times \vec{p}_{\text{THRUST}} / (|\vec{p}_\Lambda \times \vec{p}_{\text{THRUST}}|)$$

OPAL

p_T (GeV/c)	P_T^Λ (%)
< 0.3	$-1.8 \pm 3.1 \pm 1.0$
0.3 - 0.6	$0.4 \pm 1.8 \pm 0.7$
0.6 - 0.9	$1.0 \pm 1.9 \pm 0.7$
0.9 - 1.2	$0.8 \pm 2.2 \pm 0.6$
1.2 - 1.5	$0.0 \pm 2.7 \pm 0.6$
> 1.5	$1.8 \pm 1.6 \pm 0.5$
> 0.3	$0.9 \pm 0.9 \pm 0.3$
> 0.6	$1.1 \pm 1.0 \pm 0.4$

Figure 6: Measured transverse polarization of Λ baryons as a function of p_T (the transverse momentum of the Λ measured relative to the event thrust axis). The first error is statistical, second systematic.

AUTHORS ALSO FOUND NO EFFECT WITH RESPECT TO SCATTERING PLANE.

If $K^- \rightarrow \Lambda + K^0$ & $K^+ \rightarrow \bar{\Lambda} + K^0$ produce
POLARIZATION

Why do we NOT see $\Lambda \bar{\Lambda}$
TRANSVERSE POLARIZATION IN
 Z decay?

1.) ALL HADRONIC PROCESS IS UNRELATED

2.) $\pi^- \rightarrow \Lambda$ HAS OPPOSITE SIGN, AND
DATA INCLUDE Λ 'S FROM BOTH SOURCES
(d & s QUARKS).

3.) MOST OF THE DATA ($\sim 90\%$) HAVE
 $x_E < .3$, AND MAY NOT BE SENSITIVE

4.) Λ POLARIZATION IN $K^- \rightarrow \Lambda + K^0$ & $K^+ \rightarrow \bar{\Lambda} + K^0$
COULD OCCUR WHEN THE $\bar{u}$ OR u QUARK
IS STRIPPED OFF, RATHER THAN
WHEN THE $s, \bar{s}$ IS DRESSED.